

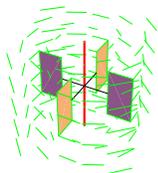
**CURL AND DIV**

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CURL (3D). The curl of a 3D vector field  $F = (P, Q, R)$  is defined as the 3D vector field

$$\text{curl}(P, Q, R) = (R_y - Q_z, P_z - R_x, Q_x - P_y).$$

CURL (2D). Recall, the curl of a 2D vector field  $F = (P, Q)$  is  $Q_x - P_y$ , a scalar field.



WANTED! Is there a multivariable calculus book, in which the above wheel is **not** shown? The wheel indicates the curl vector if  $F$  is thought of as a wind velocity field. As we will see later the direction in which the wheel turns fastest, is the direction of  $\text{curl}(F)$ . The wheel could actually be used to measure the curl of the vector field at each point. In situations with large vorticity like in a tornado, one can "see" the direction of the curl.

DIV (3D). The **divergence** of  $F = (P, Q, R)$  is the scalar field  $\text{div}(P, Q, R) = \nabla \cdot F = P_x + Q_y + R_z$ .  
 DIV (2D). The **divergence** of a vector field  $F = (P, Q)$  is  $\text{div}(P, Q) = \nabla \cdot F = P_x + Q_y$ .

NABLA CALCULUS. With the "vector"  $\nabla = (\partial_x, \partial_y, \partial_z)$ , we can write  $\text{curl}(F) = \nabla \times F$  and  $\text{div}(F) = \nabla \cdot F$ . This is both true in 2D and 3D.

LAPLACE OPERATOR.  $\Delta f = \text{divgrad}(f) = f_{xx} + f_{yy} + f_{zz}$  can be written as  $\nabla^2 f$  because  $\nabla \cdot (\nabla f) = \text{div}(\text{grad}(f))$ . One can extend this to vectorfields  $\Delta F = (\Delta P, \Delta Q, \Delta R)$  and writes  $\nabla^2 F$ .

IDENTITIES. While direct computations can verify the identities to the left, they become evident with Nabla calculus from formulas for vectors like  $\vec{v} \times \vec{v} = \vec{0}$ ,  $\vec{v} \cdot \vec{v} \times \vec{w} = 0$  or  $u \times (v \times w) = v(u \cdot w) - (u \cdot v)w$ .

$$\begin{aligned} \text{div}(\text{curl}(F)) &= 0. \\ \text{curl}(\text{grad}(F)) &= \vec{0} \\ \text{curl}(\text{curl}(F)) &= \text{grad}(\text{div}(F)) - \Delta(F). \end{aligned}$$

$$\begin{aligned} \nabla \cdot \nabla \times F &= 0. \\ \nabla \times \nabla F &= \vec{0}. \\ \nabla \times \nabla \times F &= \nabla(\nabla \cdot F) - (\nabla \cdot \nabla)F. \end{aligned}$$

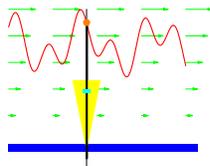
QUIZZ. Is there a vector field  $G$  such that  $F = (x + y, z, y^2) = \text{curl}(G)$ ? Answer: no because  $\text{div}(F) = 1$  is incompatible with  $\text{div}(\text{curl}(G)) = 0$ .

ADDENDA TO GREEN'S THEOREM. Green's theorem, one of the advanced topics in this course is useful in physics. We have already seen the following applications

- Simplify computation of double integrals
- Simplify computation of line integrals
- Formulas for centroid of region
- Formulas for area
- Justifying, why mechanical integrators like the planimeter works.

APROPOS PLANIMETER.

The **cone planimeter** is a mechanical instrument to find the antiderivative of a function  $f(x)$ . It uses the fact that the vector field  $F(x, y) = (y, 0)$  has  $\text{curl}(F) = -1$ . By Greens theorem, the line integral around the type I region bounded by 0 and the graph of  $f(x)$  in the counter clockwise direction is  $\int_a^x f(x) dx$ . The planimeter determines that line integral.



THERMODYNAMICS. Gases or liquids are often described in a  $P - V$  **diagram**, where the volume in the  $x$ -axis and the pressure in the  $y$  axis. A periodic process like the pump in a refrigerator leads to closed curve  $\gamma : r(t) = (V(t), p(t))$  in the  $V - p$  plane. The curve is parameterized by the time  $t$ . At a given time, the gas has volume  $V(t)$  and a pressure  $p(t)$ . Consider the vector field  $F(V, p) = (p, 0)$  and a closed curve  $\gamma$ . What is  $\int_{\gamma} F ds$ ? Writing it out, we get  $\int_0^t (p(t), 0) \cdot (V'(t), p'(t)) dt = \int_0^t p(t)V'(t)dt = \int_0^t p dV$ . The curl of  $F(V, p)$  is  $-1$ . You see by Green's theorem the integral  $-\int_0^t p dV$  is the area of the region enclosed by the curve.

MAXWELL EQUATIONS (in homework you assume no current  $j = 0$  and charges  $\rho = 0$ .  $c =$  speed of light.)

$\text{div}(B) = 0$	No monopoles	there are no magnetic monopoles.
$\text{curl}(E) = -\frac{1}{c}B_t$	Faraday's law	change of magnetic flux induces voltage
$\text{curl}(B) = \frac{1}{c}E_t + \frac{4\pi}{c}j$	Ampère's law	current or change of E produces magnetic field
$\text{div}(E) = 4\pi\rho$	Gauss law	electric charges are sources for electric field

2D MAXWELL EQUATIONS? In space dimensions  $d$  different than 3 the electromagnetic field has  $d(d+1)/2$  components. In 2D, the magnetic field is a scalar field and the electric field  $E = (P, Q)$  a vector field. The 2D Maxwell equations are  $\text{curl}(E) = -\frac{1}{c}\frac{d}{dt}B$ ,  $\text{div}(E) = 4\pi\rho$ . Consider a region  $R$  bounded by a wire  $\gamma$ . Green's theorem tells us that  $d/dt \int_R B(t) dx dy$  is the line integral of  $E$  around the boundary. But  $\int_{\gamma} E ds$  is a voltage. A change of the magnetic field produces a voltage. This is the **dynamo** in 2D. We will see the real dynamo next week in 3D, where electromagnetism works (it would be difficult to generate a magnetic field in flatland). If  $v$  is the velocity distribution of a fluid in the plane, then  $\omega(x, y) = \text{curl}(v(x, y))$  is the **vorticity** of the fluid. The integral  $\int_R \omega dx dy$  is called the **vortex flux** through  $R$ . Green's theorem assures that this flux is related to the circulation on the boundary.

FLUID DYNAMICS.  $v$  velocity,  $\rho$  density of fluid.

Continuity equation	$\dot{\rho} + \text{div}(\rho v) = 0$	no fluid get lost
Incompressibility	$\text{div}(v) = 0$	incompressible fluids, in 2D: $v = \text{grad}(u)$
Irrotational	$\text{curl}(v) = 0$	irrotation fluids

A RELATED THEOREM. If we rotate the vector field  $F = (P, Q)$  by 90 degrees  $= \pi/2$  we get a new vector field  $G = (-Q, P)$ . The integral  $\int_{\gamma} F \cdot ds$  becomes a **flux**  $\int_{\gamma} G \cdot dn$  of  $G$  through the boundary of  $R$ , where  $dn$  is a normal vector with the length of  $dr$ . With  $\text{div}(F) = (P_x + Q_y)$ , we see that  $\text{curl}(F) = \text{div}(G)$ . Green's theorem now becomes

$$\iint_R \text{div}(G) dx dy = \int_{\gamma} G \cdot dn,$$

where  $dn(x, y)$  is a normal vector at  $(x, y)$  orthogonal to the velocity vector  $r'(x, y)$  at  $(x, y)$ . This new theorem has a generalization to three dimensions, where it is called Gauss theorem or divergence theorem. Don't treat this however as a different theorem in two dimensions. **It is in two dimensions just Green's theorem disguised.** There are only 2 basic integral theorems in the plane: Green's theorem and the FTLI.

PREVIEW. Green's theorem is of the form  $\int_R F' = \int_{\delta R} F$ , where  $F'$  is a "derivative" and  $\delta R$  is a "boundary". There are 2 such theorems in dimensions 2, three theorems in dimensions 3, four in dimension 4 etc. In the plane, Green's theorem is the second one besides the fundamental theorem of line integrals FTLI. In three dimensions, there are two more theorems beside the FTLI: Stokes and Gauss Theorems which we will see in the next week.

dim	dim(R)	theorem
1D	1	Fund. thm of calculus

dim	dim(R)	theorem
3D	1	Fundam. thm of line integrals
3D	2	Stokes theorem
3D	3	Gauss theorem

2D	1	Fund. thm of line integrals
2D	2	Green's theorem

$1 \mapsto 1$	$f'$	derivative
$1 \mapsto 2$	$\nabla f$	gradient
$2 \mapsto 1$	$\nabla \times F$	curl

$1 \mapsto 3$	$\nabla f$	gradient
$3 \mapsto 3$	$\nabla \times F$	curl
$3 \mapsto 1$	$\nabla \cdot F$	divergence