

## FOURIER SERIES

Math 21b, Spring 07

Smooth functions  $f(x)$  on  $[-\pi, \pi]$  form a linear space  $X$ . There is an **inner product** in  $X$  defined by

$$\langle f, g \rangle = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x)g(x) dx$$

It allows to define angles, length, projections etc in the space  $X$  as we did in finite dimensions.

THE FOURIER BASIS.

**THEOREM.** The functions  $\{\cos(nx), \sin(nx), 1/\sqrt{2}\}$  form an orthonormal basis in  $X$ .

Proof. To check linear independence we have to evaluate a few integrals. For all  $n, m \geq 1$ , with  $n \neq m$  we have

$$\begin{aligned} \langle \frac{1}{\sqrt{2}}, \frac{1}{\sqrt{2}} \rangle &= 1 \\ \langle \cos(nx), \cos(nx) \rangle &= 1, \langle \cos(nx), \cos(mx) \rangle = 0 \\ \langle \sin(nx), \sin(nx) \rangle &= 1, \langle \sin(nx), \sin(mx) \rangle = 0 \\ \langle \sin(nx), \cos(mx) \rangle &= 0 \\ \langle \sin(nx), 1/\sqrt{2} \rangle &= 0 \\ \langle \cos(nx), 1/\sqrt{2} \rangle &= 0 \end{aligned}$$

To verify the above integrals in the homework, the following trigonometric identities are useful:

$$\begin{aligned} 2 \cos(nx) \cos(my) &= \cos(nx - my) + \cos(nx + my) \\ 2 \sin(nx) \sin(my) &= \cos(nx - my) - \cos(nx + my) \\ 2 \sin(nx) \cos(my) &= \sin(nx + my) + \sin(nx - my) \end{aligned}$$

FOURIER COEFFICIENTS. The **Fourier coefficients** of  $f$  are defined as

$$a_0 = \langle f, 1/\sqrt{2} \rangle = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x)/\sqrt{2} dx$$

$$a_n = \langle f, \cos(nt) \rangle = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos(nx) dx$$

$$b_n = \langle f, \sin(nt) \rangle = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx) dx$$

FOURIER SERIES. The **Fourier representation** of a smooth function  $f$  is the identity

$$f(x) = \frac{a_0}{\sqrt{2}} + \sum_{k=1}^{\infty} a_k \cos(kx) + \sum_{k=1}^{\infty} b_k \sin(kx)$$

We take it for granted that the series converges and that the identity holds for all  $x$ .

ODD AND EVEN FUNCTIONS. Here is some advice which can save time when computing Fourier series:

If  $f$  is odd:  $f(x) = -f(-x)$ , then  $f$  has a sin series.

If  $f$  is even:  $f(x) = f(-x)$ , then  $f$  has a cos series.

If you integrate an odd function over  $[-\pi, \pi]$  you get 0.

The product of two odd functions is even, the product between an even and an odd function is odd.

EXAMPLE 1. Let  $f(x) = x$  on  $[-\pi, \pi]$ . This is an odd function ( $f(-x) + f(x) = 0$ ) so that it has a sin series: with  $b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} x \sin(nx) dx = \frac{-1}{\pi} (x \cos(nx)/n + \sin(nx)/n^2) \Big|_{-\pi}^{\pi} = 2(-1)^{n+1}/n$ , we get  $x = \sum_{n=1}^{\infty} 2 \frac{(-1)^{n+1}}{n} \sin(nx)$ . For example

$$\frac{\pi}{2} = 2 \left( \frac{1}{1} - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \dots \right)$$

is a **formula of Leibnitz**.

EXAMPLE 2. Let  $f(x) = \cos(x) + 1/7 \cos(5x)$ . This **trigonometric polynomial** is already the Fourier series. The nonzero coefficients are  $a_1 = 1, a_5 = 1/7$ .

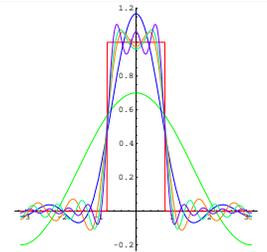
EXAMPLE 3. Let  $f(x) = 1$  on  $[-\pi/2, \pi/2]$  and  $f(x) = 0$  else. This is an even function  $f(-x) - f(x) = 0$  so that it has a cos series: with  $a_0 = 1/(\sqrt{2}), a_n = \frac{1}{\pi} \int_{-\pi/2}^{\pi/2} 1 \cos(nx) dx = \frac{\sin(nx)}{\pi n} \Big|_{-\pi/2}^{\pi/2} = \frac{2(-1)^m}{\pi(2m+1)}$  if  $n = 2m + 1$  is odd and 0 else. So, the series is

$$f(x) = \frac{1}{2} + \frac{2}{\pi} \left( \frac{\cos(x)}{1} - \frac{\cos(3x)}{3} + \frac{\cos(5x)}{5} - \dots \right)$$

**Remark.** The function in Example 3 is not smooth but Fourier theory still works. What happened at the discontinuity point  $\pi/2$ ? The Fourier series gives 0. Diplomatically, it has chosen the point in the middle of the limits from the right and the limit from the left.

FOURIER APPROXIMATION. For a smooth function  $f$ , the Fourier series of  $f$  converges to  $f$ . The Fourier coefficients are the coordinates of  $f$  in the Fourier basis.

The function  $f_n(x) = \sum_{k=1}^n a_k \sin(kx)$  is called a **Fourier approximation** of  $f$ . The picture to the right shows an approximation of a piecewise continuous even function in example 3).



THE PARSEVAL EQUALITY. When evaluating the square of the length of  $f$  with the square of the length of the series, we get

$$\|f\|^2 = a_0^2 + \sum_{k=1}^{\infty} a_k^2 + b_k^2$$

EXAMPLE. We have seen in example 1 that  $f(x) = x = 2(\sin(x) - \sin(2x)/2 + \sin(3x)/3 - \sin(4x)/4 + \dots)$ . Because the Fourier coefficients are  $a_k = 2(-1)^{k+1}/k$ , we have  $4(1 + 1/4 + 1/9 + \dots) = \frac{1}{\pi} \int_{-\pi}^{\pi} x^2 dx = 2\pi^2/3$  and so

$$\frac{1}{1} + \frac{1}{4} + \frac{1}{9} + \frac{1}{16} + \frac{1}{25} + \dots = \frac{\pi^2}{6}$$

Isn't it fantastic that we can sum up the reciprocal squares? This formula has been obtained already by Leonard Euler, who celebrated his 300th birthday two weeks ago, on April 15.

HOMEWORK: (this homework is due either Wednesday or Thursday. Remember that there is a lab due on Thursday or Friday).

1. Verify that the functions  $\cos(nx), \sin(nx), 1/\sqrt{2}$  form an orthonormal family.
2. Find the Fourier series of the function  $f(x) = |x|$ .
3. Find the Fourier series of the function  $\cos^2(x) + 5 \sin(x) + 5$ . You may find the double angle formula  $\cos^2(x) = \frac{\cos(2x)+1}{2}$  useful.
4. Find the Fourier series of the function  $f(x) = |\sin(x)|$ .
5. In the previous problem 4) you should have gotten a series

$$f(x) = \frac{2}{\pi} - \frac{4}{\pi} \left( \frac{\cos(2x)}{2^2-1} + \frac{\cos(4x)}{4^2-1} + \frac{\cos(6x)}{6^2-1} + \dots \right)$$

Use Parseval's identity to find the value of

$$\frac{1}{(2^2-1)^2} + \frac{1}{(4^2-1)^2} + \frac{1}{(6^2-1)^2} + \dots$$

## HEAT AND WAVE EQUATION

Math 21b, Spring 07

**FUNCTIONS OF TWO VARIABLES.** We consider functions  $f(x, t)$  which are for fixed  $t$  a piecewise smooth function in  $x$ . Analogously as we studied the motion of a **vector**  $\vec{v}(t)$ , we are now interested in the motion of a **function**  $f$  in time  $t$ . While the governing equation for a vector was an ordinary differential equation  $\dot{x} = Ax$  (ODE), the describing equation is now be a **partial differential equation** (PDE)  $\dot{f} = T(f)$ . The function  $f(x, t)$  could denote the **temperature of a stick** at a position  $x$  at time  $t$  or the **displacement of a string** at the position  $x$  at time  $t$ . The motion of these dynamical systems will be easy to describe in the orthonormal Fourier basis  $1/\sqrt{2}, \sin(nx), \cos(nx)$  treated in an earlier lecture.

**PARTIAL DERIVATIVES.** We write  $f_x(x, t)$  and  $f_t(x, t)$  for the **partial derivatives** with respect to  $x$  or  $t$ . The notation  $f_{xx}(x, t)$  means that we differentiate twice with respect to  $x$ .

Example: for  $f(x, t) = \cos(x + 4t^2)$ , we have

- $f_x(x, t) = -\sin(x + 4t^2)$
- $f_t(x, t) = -8t \sin(x + 4t^2)$ .
- $f_{xx}(x, t) = -\cos(x + 4t^2)$ .

One also uses the notation  $\frac{\partial f(x, y)}{\partial x}$  for the partial derivative with respect to  $x$ . Tired of all the "partial derivative signs", we always write  $f_x(x, t)$  for the partial derivative with respect to  $x$  and  $f_t(x, t)$  for the partial derivative with respect to  $t$ .

**PARTIAL DIFFERENTIAL EQUATIONS.** A partial differential equation is an equation for an unknown function  $f(x, t)$  in which different partial derivatives occur.

- $f_t(x, t) + f_x(x, t) = 0$  with  $f(x, 0) = \sin(x)$  has a solution  $f(x, t) = \sin(x - t)$ .
- $f_{tt}(x, t) - f_{xx}(x, t) = 0$  with  $f(x, 0) = \sin(x)$  and  $f_t(x, 0) = 0$  has a solution  $f(x, t) = (\sin(x - t) + \sin(x + t))/2$ .

**THE HEAT EQUATION.** The temperature distribution  $f(x, t)$  in a metal bar  $[0, \pi]$  satisfies the **heat equation**

$$f_t(x, t) = \mu f_{xx}(x, t)$$

This partial differential equation tells that the rate of change of the temperature at  $x$  is proportional to the second space derivative of  $f(x, t)$  at  $x$ . The function  $f(x, t)$  is assumed to be zero at both ends of the bar and  $f(x) = f(x, 0)$  is a given initial temperature distribution. The constant  $\mu$  depends on the heat conductivity properties of the material. Metals for example conduct heat well and would lead to a large  $\mu$ .

**REWRITING THE PROBLEM.** We can write the problem as

$$\frac{d}{dt} f = \mu D^2 f$$

We will solve the problem in the same way as we solved linear differential equations:

$$\frac{d}{dt} \vec{x} = A \vec{x}$$

where  $A$  is a matrix - by diagonalization.

We use that the Fourier basis is just the diagonalization:  $D^2 \cos(nx) = -n^2 \cos(nx)$  and  $D^2 \sin(nx) = -n^2 \sin(nx)$  show that  $\cos(nx)$  and  $\sin(nx)$  are eigenfunctions to  $D^2$  with eigenvalue  $n^2$ . By a symmetry trick, we can focus on sin-series from now on.

**SOLVING THE HEAT EQUATION WITH FOURIER THEORY.** The heat equation  $f_t(x, t) = \mu f_{xx}(x, t)$  with smooth  $f(x, 0) = f(x)$ ,  $f(0, 0) = f(\pi, 0) = 0$  has the solution

$$f(x, t) = \sum_{n=1}^{\infty} b_n \sin(nx) e^{-n^2 \mu t}$$

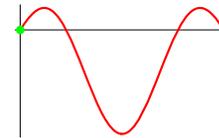
$$b_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$$

Proof: With the initial condition  $f(x) = \sin(nx)$ , we have the evolution  $f(x, t) = e^{-\mu n^2 t} \sin(nx)$ . If  $f(x) = \sum_{n=1}^{\infty} b_n \sin(nx)$  then  $f(x, t) = \sum_{n=1}^{\infty} b_n e^{-\mu n^2 t} \sin(nx)$ .

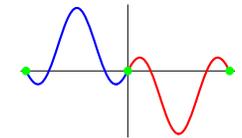
**A SYMMETRY TRICK.** Given a function  $f$  on the interval  $[0, \pi]$  which is zero at 0 and  $\pi$ . It can be extended to an odd function on the doubled interval  $[-\pi, \pi]$ .

The Fourier series of an odd function is a pure sin-series. The Fourier coefficients are  $b_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$ .

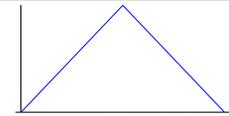
The function is given on  $[0, \pi]$ .



The odd symmetric extension on  $[-\pi, \pi]$ .



**EXAMPLE.** Assume the initial temperature distribution  $f(x, 0)$  is a sawtooth function which has slope  $\pi/2$  on the interval  $[0, \pi/2]$  and slope  $-\pi/2$  on the interval  $[\pi/2, \pi]$ . We first compute the sin-Fourier coefficients of this function.

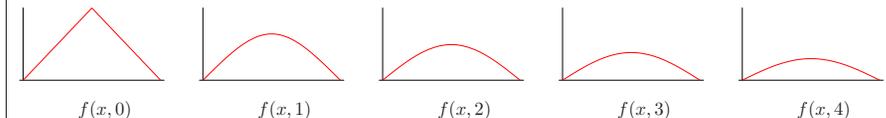


The sin-Fourier coefficients are  $b_n = \frac{4}{n^2 \pi} (-1)^{(n-1)/2}$  for odd  $n$  and 0 for even  $n$ . The solution is

$$f(x, t) = \sum_n b_n e^{-\mu n^2 t} \sin(nx)$$

The exponential term containing the time makes the function  $f(x, t)$  converge to 0: The body cools. The higher frequencies are damped faster: "smaller disturbances are smoothed out faster."

**VISUALIZATION.** We can plot the graph of the function  $f(x, t)$  or slice this graph and plot the temperature distribution for different values of the time  $t$ .



**THE WAVE EQUATION.** The height of a string  $f(x, t)$  at time  $t$  and position  $x$  on  $[0, \pi]$  satisfies the **wave equation**

$$f_{tt}(t, x) = c^2 f_{xx}(t, x)$$

where  $c$  is a constant. As we will see,  $c$  is the **speed** of the waves.

REWRITING THE PROBLEM. We can write the problem as

$$\frac{d^2}{dx^2} f = c^2 D^2 f$$

We will solve the problem in the same way as we solved

$$\frac{d^2}{dx^2} \vec{x} = A\vec{x}$$

If  $A$  is diagonal, then every basis vector  $x$  satisfies an equation of the form  $\frac{d^2}{dx^2} x = -c^2 x$  which has the solution  $x(t) = x(0) \cos(ct) + x'(0) \sin(ct)/c$ .

SOLVING THE WAVE EQUATION WITH FOURIER THEORY. The wave equation  $f_{tt} = c^2 f_{xx}$  with  $f(x, 0) = f(x), f_t(x, 0) = g(x), f(0, t) = f(\pi, t) = 0$  has the solution

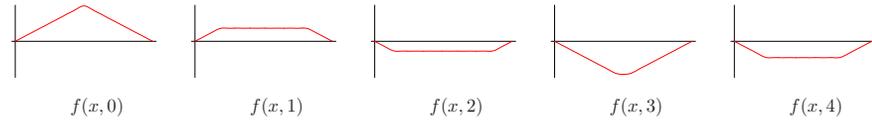
$$f(x, t) = \sum_{n=1}^{\infty} \frac{a_n}{nc} \sin(nx) \cos(nct) + \sum_{n=1}^{\infty} \frac{b_n}{nc} \sin(nx) \sin(nct)$$

$$a_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$$

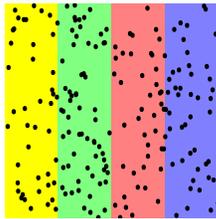
$$b_n = \frac{2}{\pi} \int_0^{\pi} g(x) \sin(nx) dx$$

Proof. With  $f(x) = \sin(nx), g(x) = 0$ , the solution is  $f(x, t) = \cos(ct) \sin(nx)$ . With  $f(x) = 0, g(x) = \sin(nx)$ , the solution is  $f(x, t) = \frac{1}{c} \sin(ct) \sin(nx)$ . For  $f(x) = \sum_{n=1}^{\infty} a_n \sin(nx)$  and  $g(x) = \sum_{n=1}^{\infty} b_n \sin(nx)$ , we get the formula by summing these two solutions.

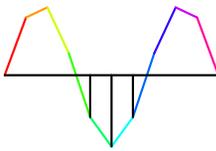
VISUALIZATION. We can just plot the graph of the function  $f(x, t)$  or plot the string for different times  $t$ .



TO THE DERIVATION OF THE HEAT EQUATION. The temperature  $f(x, t)$  is proportional to the kinetic energy at  $x$ . Divide the stick into  $n$  adjacent cells and assume that in each time step, a fraction of the particles moves randomly either to the right or to the left. If  $f_i(t)$  is the **energy** of particles in cell  $i$  at time  $t$ , then the energy of particles at time  $t + 1$  is proportional to  $(f_{i-1}(t) - 2f_i(t) + f_{i+1}(t))$ . This is a discrete version of the second derivative because  $dx^2 f_{xx}(t, x) \sim (f(x + dx, t) - 2f(x, t) + f(x - dx, t))$ .



TO THE DERIVATION OF THE WAVE EQUATION. We can model a string by  $n$  discrete particles linked by strings. Assume that the particles can move up and down only. If  $f_i(t)$  is the **height** of the particles, then the right particle pulls with a force  $f_{i+1} - f_i$ , the left particle with a force  $f_{i-1} - f_i$ . Again,  $(f_{i-1}(t) - 2f_i(t) + f_{i+1}(t))$  which is a discrete version of the second derivative because  $dx^2 f_{xx}(t, x) \sim (f(x + dx, t) - 2f(x, t) + f(x - dx, t))$ .



OVERVIEW: The heat and wave equation can be solved like ordinary differential equations:

<p>Ordinary differential equations</p> $x_t(t) = Ax(t)$ $x_{tt}(t) = Ax(t)$	<p>Partial differential equations</p> $f_t(t, x) = f_{xx}(t, x)$ $f_{tt}(t, x) = f_{xx}(t, x)$
<p>Diagonalizing <math>A</math> leads for eigenvectors <math>\vec{v}</math></p> $Av = -c^2 v$	<p>Diagonalizing <math>T = D^2</math> with eigenfunctions <math>f(x) = \sin(nx)</math></p> $Tf = -n^2 f$
<p>to the differential equations</p> $v_t = -c^2 v$ $v_{tt} = -c^2 v$	<p>leads to the differential equations</p> $f_t(x, t) = -n^2 f(x, t)$ $f_{tt}(x, t) = -n^2 f(x, t)$
<p>which are solved by</p> $v(t) = e^{-c^2 t} v(0)$ $v(t) = v(0) \cos(ct) + v_t(0) \sin(ct)/c$	<p>which are solved by</p> $f(x, t) = f(x, 0) e^{-n^2 t}$ $f(x, t) = f(x, 0) \cos(nt) + f_t(x, 0) \sin(nt)/n$

NOTATION:

$f$  function on  $[-\pi, \pi]$  smooth or piecewise smooth.  $Tf = \lambda f$  Eigenvalue equation analogously to  $Av = \lambda v$ .  
 $t$  time variable  $f_t$  partial derivative of  $f(x, t)$  with respect to time  $t$ .  
 $x$  space variable  $f_x$  partial derivative of  $f(x, t)$  with respect to space  $x$ .  
 $D$  the partial differential operator  $Df(x) = f'(x) = f_{xx}$  second partial derivative of  $f$  twice with respect to space  $x$ .  
 $d/dx f(x)$ .  
 $T$  linear transformation, like  $Tf = D^2 f = f''$ .  
 $c$  speed of the wave.  $\mu$  heat conductivity  
 $f(x) = -f(-x)$  odd function, has sin Fourier series

HOMEWORK. This homework is due until Tuesday morning May 7 in the mailboxes of your CA:

1) Solve the heat equation  $f_t = \mu f_{xx}$  on  $[0, \pi]$  with the initial condition  $f(x, 0) = |\sin(3x)|$ .

We want to see in exercises 2-4 how to deal with solutions to the heat equation, where the boundary values are not 0. Problems 2-4 belong together:

2) Verify that for any constants  $a, b$  the function  $h(x, t) = (b - a)x/\pi + a$  is a solution to the heat equation.

3) Assume we have the problem to describe solutions  $f(x, t)$  to the heat equations, where  $f(0, t) = a$  and  $f(\pi, t) = b$ . Show that  $f(x, t) - h(x, t)$  is a solution of the heat equation with  $f(0, t) = 0$  and  $f(\pi, t) = 0$ .

4) Solve the heat equation with the initial condition  $f(x, 0) = f(x) = \sin(3x) + x/\pi$  and satisfying  $f(0, t) = 0, f(\pi, t) = 1$  for all times  $t$ . This is a situation, when the stick is kept at constant but different temperatures on the both ends.

5) A piano string is fixed at the ends  $x = 0$  and  $x = \pi$  and initially undisturbed. The piano hammer induces an initial velocity  $u_t(x, t) = g(x)$  onto the string, where  $g(x) = \sin(2x)$  on the interval  $[0, \pi/2]$  and  $g(x) = 0$  on  $[\pi/2, \pi]$ . Find the motion of the string.